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PNO: a promising deep-UV nonlinear optical material with the largest second harmonic generation effect†

 Congwei Xie,[‡]*^a Abudukadi Tudi[‡]^b and Artem R. Oganov[‡]*^a

In this work, the polar tetrahedron [PN₂O₂] was revealed as a new deep-ultraviolet (deep-UV) nonlinear optically active unit. Accordingly, a thermodynamically stable compound (PNO) consisting of the polar [PN₂O₂] units was predicted and suggested as a promising candidate for deep-UV nonlinear optical (NLO) materials. Compared with other deep-UV materials known to date, PNO possesses the strongest second harmonic generation (SHG) coefficient (about 6 times that of KH₂PO₄ (KDP)). Moreover, its three-dimensional connectivity endows it with good mechanical and thermal properties. Therefore, PNO should be a new option for non- π -conjugated deep-UV NLO materials.

Nonlinear optical (NLO) crystals for producing coherent light in the deep-ultraviolet (deep-UV, $\lambda \leq 200$ nm) range are the key materials of solid-state deep-UV lasers.¹ To date, KBe₂BO₃F₂ (KBBF) is still the only practical deep-UV NLO crystal which can directly generate lasers with wavelengths shorter than 200 nm by second harmonic generation (SHG), although its further industrial applications are seriously limited by two drawbacks: it contains highly toxic beryllium and it is hard to grow its large crystals because of the platelet-like shape of crystals.² Therefore, exploring new deep-UV NLO materials is urgently needed and of great interest from both scientific and industrial standpoints.^{3–6}

A promising deep-UV NLO material should simultaneously satisfy many stringent criteria, such as high deep-UV transparency (band gap $E_g \geq 6.2$ eV), relatively large SHG coefficient (at least comparable to that of the classical NLO crystal KH₂PO₄

(KDP) which is 0.39 pm V⁻¹), appropriate birefringence for phase-matching (PM) and so on. For the past few decades, materials that could meet these conditions were mainly found in π -conjugated systems, *e.g.*, borates, carbonates and their derivatives, in which the π -conjugated building units ([BO₃], [B₃O₆], [CO₃], *etc.*) are beneficial to the increase of both the SHG coefficient and birefringence.⁷

In addition to π -conjugated systems, non- π -conjugated systems such as silicates, phosphates and sulfates have also been considered as potential candidates for deep-UV NLO materials, where tetrahedral groups [XO₄] (X = P, S, Si) are favorable with deep-UV transparency.⁸ Many NLO-active phosphates, sulfates and silicates with high deep-UV transparency and SHG coefficient have been discovered. Unfortunately, none of them possess sufficient birefringence to satisfy the PM conditions in the deep-UV region because of the weak optical anisotropy of the nonpolar [XO₄] tetrahedral groups in their structures. Thus, regarding non- π -conjugated NLO materials, the greatest challenge lies in how to achieve enough birefringence to realize PM in the deep-UV region.⁹

The common strategy to improve birefringence is to introduce transition metal cations with d^0 or d^{10} electronic configuration or ns^2 lone pair cations.^{10,11} However, in most cases, this causes a heavy red-shift of the absorption edge and thus does not work to design deep-UV materials. Recently, polar polyhedra in which the central positive ions are bonded to two anionic ligands (*e.g.*, O, F, NH₂, *etc.*) were proposed as good functional structural units (FSUs) for designing deep-UV NLO materials.^{3,12–14} Among them, polar tetrahedra [BO_{4-x}F_x], [PO_{4-x}F_x], [SO_{4-x}F_x] and [SO_{4-x}(NH₂)_x] ($x = 1, 2$) have been found to possess large polarizability anisotropy and hyperpolarizability and thus have advantages on enhancing birefringence and SHG coefficients.^{3,12–14} Accordingly, several fluoroborates, fluorophosphates and sulfamides have been theoretically reported and/or experimentally confirmed as deep-UV NLO materials.^{3,12b,13a,14}

We noticed that nitrogen can also be used to substitute oxygen in [XO₄] units. It has been reported that the introduction of nitrogen into a [SiO₄] unit can form polar [SiNO₃],

^a Skolkovo Institute of Science and Technology, Skolkovo Innovation Center, Moscow 121205, Russian Federation. E-mail: cwxie@ms.xjtu.ac.cn, a.oganov@skoltech.ru

^b CAS Key Laboratory of Functional Materials and Devices for Special Environments, Xinjiang Technical Institute of Physics & Chemistry, CAS, Xinjiang Key Laboratory of Electronic Information Materials and Devices, 40-1 South Beijing Road, Urumqi 830011, China

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‡ These authors have contributed to this work equally.

[SiN₂O₂], and [SiN₃O] tetrahedra with large polarizability anisotropy and high hyperpolarizability, and oxonitridosilicates containing these polar tetrahedra have been identified as new candidate systems to explore new UV NLO materials.¹⁵

In this work, our theoretical calculations showed that the polar [PN₂O₂] tetrahedron has large polarizability anisotropy and hyperpolarizability, and relatively large highest occupied molecular orbital (HOMO)–lowest unoccupied molecular orbital (LUMO) gap and optical gap, indicating that the [PN₂O₂] tetrahedron should be a good deep-UV NLO-active unit. To confirm this, a thermodynamically stable noncentrosymmetric (NCS) compound (PNO) consisting of polar [PN₂O₂] tetrahedra was predicted for evaluating its optical properties. The present theoretical study showed that PNO has a short deep-UV cut-off edge of about 188 nm and a PM wavelength less than 190 nm. In addition, benefiting from both the excellent intrinsic optical properties of [PN₂O₂] tetrahedra and their ordered arrangement, PNO achieves the strongest SHG coefficient (about 6 times that of KDP) among all deep-UV NLO materials. Similar to [PN₂O₂] tetrahedra, [PNO₃] is also newly suggested to be a potential DUV-NLO unit. However, we did not predict any stable ternary P_xN_yO_z compounds containing [PO₃N] tetrahedra.

We have evaluated the basic electronic and optical properties for a series of Si and P based tetrahedral anionic groups: [SiO_{4-x}F_x] ($x = 0, 1, 2, 3$), [SiN_xO_{4-x}] ($x = 1, 2, 3, 4$), [PO_{4-x}F_x] ($x = 0, 1, 2$) and [PN_xO_{4-x}] ($x = 1, 2, 3, 4$). We first performed geometry optimization and then investigated their electronic structures and optical properties based on DFT calculations. As shown in Fig. 1 and listed in Table S2 in the ESI[†], the polar [PN₂O₂] and [PNO₃] units exhibit high polarizability anisotropy values of 35.6 and 11.9, respectively. As a comparison, the polarizability anisotropy values of the well-known deep-UV NLO-active units [PO₂F₂] and [PO₃F] calculated by using the same computational settings are 6.7 and 6.0, respectively. The [PN₂O₂] and [PNO₃] units also have larger hyperpolarizabilities than [PO₂F₂] and [PO₃F] units. Although the substitution of O by N will reduce the HOMO–LUMO gap and optical gap, [PN₂O₂] and [PNO₃] units still have relatively large HOMO–LUMO gaps

(> 5 eV) and optical gaps (> 4 eV). We here note that, in P_xN_yO_z compounds, N atoms will be coordinated with more than one P atom, which can eliminate the non-bonding electrons and thus helpful to enhance the band gap.^{15b} Thus, similar to [PO₃F] and [PO₂F₂] units, [PN₂O₂] and [PNO₃] units can also balance the HOMO–LUMO gap, hyperpolarizabilities and polarizability anisotropies, suggesting that they are good deep-UV NLO active units.

To explore potential deep-UV P_xN_yO_z compounds, using an evolutionary algorithm as implemented in the USPEX code^{16–18}, we performed variable composition structure searches for the P–N–O system with up to 36 atoms in the unit cell at zero temperature and ambient pressure. Based on the total energy computed for all the predicted structures, we constructed the phase diagram for the P–N–O system. As shown in Fig. 2, along with the experimentally reported P₄N₆O¹⁹, we can find that PNO is also thermodynamically stable. Luckily, the lowest-energy PNO structure (**PNO-I**) is a non-centrosymmetric structure (see Table S3 in the ESI[†]). In addition to the stable non-centrosymmetric structure, we have also predicted one competitive metastable non-centrosymmetric structure (**PNO-II**); its energy is slightly higher (by 1 meV per atom) than that of **PNO-I**. We here note that both **PNO-I** and **PNO-II** can be also produced by the structural analogy technique, as reported by Materials Project.²⁰ For **PNO-II**, along with other metastable PNO structures predicted using USPEX, details of their structural and thermodynamic information are also listed in Table S3 in the ESI[†].

With fully relaxed structures, we computed phonon dispersion curves and elastic constants for confirming the mechanical and dynamical stability of **PNO-I** and **PNO-II**. It was found that both these PNO structures are mechanically and dynamically stable (see Fig. S2 and Table S4 in the ESI[†]). The crystal structures of **PNO-I** and **PNO-II** are shown in Fig. 3. The asymmetric unit of **PNO-I** contains one unique P atom, one unique O atom and one unique N atom and each P atom is connected to two O atoms and two N atoms forming the FSU [PN₂O₂] tetrahedra (see Fig. 3(a)). **PNO-I** belongs to the *I*₂12₁2₁ space group (see Fig. 3(b)), and the corner-sharing [PN₂O₂]

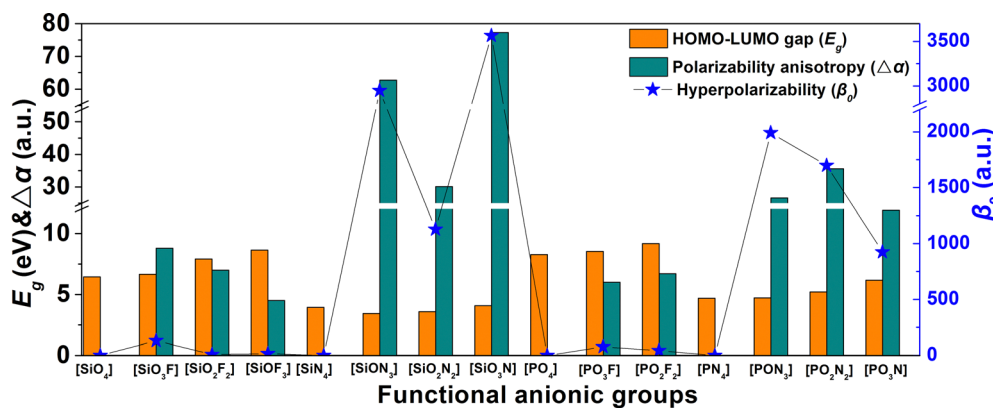


Fig. 1 Computed HOMO–LUMO gaps (E_g , in eV), polarizability anisotropies (δ , in a.u., 1 a.u. = $1.6488 \times 10^{-41} \text{ C}^2 \text{ m}^2 \text{ J}^{-1}$) and the mean static first hyperpolarizability (β_0 , in a.u., 1 a.u. = $3.20636 \times 10^{-53} \text{ C}^3 \text{ m}^3 \text{ J}^{-2}$) of [SiO_{4-x}F_x] ($x = 0, 1, 2, 3$), [SiO_{4-x}N_x] ($x = 1, 2, 3, 4$), [PO_{4-x}F_x] ($x = 0, 1, 2$) and [PO_{4-x}N_x] ($x = 1, 2, 3, 4$) tetrahedral anionic groups.

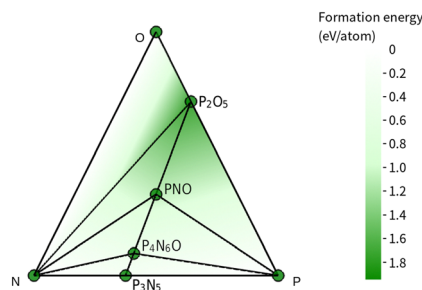


Fig. 2 Phase diagram of the P–N–O system. Only stable compounds are shown and they are denoted by circles.

tetrahedra form a three-dimensional framework structure (see Fig. 3(c)), derived from tridymite (see Fig. 3(d)). **PNO-II** also has one unique P atom, one unique O atom and one unique N atom and has the same FSU found in **PNO-I**. **PNO-II** belongs to the *Cc* space group (see Fig. 3(e)) and also has a 3D framework structure derived from cristobalite (see Fig. 3(f) and (g)). We here note that this 3D-connectivity nature of PNO is helpful to yield superior properties such as good mechanical strength and thermal conductivity (see Table S4 in the ESI†) and thus may be beneficial to its further applications.

The NLO properties of **PNO-I** and **PNO-II** are shown in Table 1. Both predicted structures have relatively wide band gaps greater than 6.2 eV, which demonstrates their ability to achieve deep-UV transmission. They both have large birefringence ($>0.151@1064$ nm) beneficial to realize PM. The PM wavelengths of **PNO-I** and **PNO-II** are lower than 190 nm (see Fig. S3 in the ESI†), and the static SHG coefficients of these two structures are large. The SHG coefficients and PM wavelengths of PNO structures along with those of other well-known deep-UV materials are shown in Fig. S5 (ESI†). We can find that **PNO-I** and **PNO-II** have very large SHG coefficients (about 6–9 times that of KDP). Compared with other compounds containing non- π -conjugated groups and π -conjugated groups, PNO has the largest SHG coefficient.²¹

In order to resolve the source of the excellent NLO properties of PNO, we performed structural performance analysis using first-principles calculations for **PNO-I**. As shown in Fig. 4(a), **PNO-I** has direct band gap with an HSE06-value of 6.61 eV. The

Table 1 Calculated band gaps using the hybrid HSE06 functional (E_g -HSE, in eV), the independent SHG coefficients at zero frequency (SHG, in pm V^{-1}), birefringence ($\Delta n@1064$ nm), and the shortest phase-matching wavelengths (λ_{PM} , in nm) for **PNO-I** and **PNO-II**

Compound	Space groups	E_g -HSE	SHG	Δn	λ_{PM}
PNO-I	Cal $I2_12_12_1$	6.61	$d_{36} = 2.33$	0.151	<190
PNO-II	Cal <i>Cc</i>	6.63	$d_{11} = -3.44$; $d_{31} = 0.89$ $d_{26} = 2.16$; $d_{35} = 0.68$ $d_{32} = 1.01$; $d_{33} = -0.87$	0.174	<190

valence bands (VBs) are mainly composed of N-2p states, O-2p states and P-3p states, and the conduction bands (CBs) are mainly derived from P-3p states (Fig. 4(b)). The SHG density and band-resolved methods²² were used to analyze the largest SHG coefficient of **PNO-I** (Fig. 4(b)–(d) and Fig. S6 in the ESI†).

The virtual-electron (VE) and the virtual-hole (VH) transitions account for 72% and 28% of the SHG effect, respectively. The VE process is shown in Fig. 4(c) and (d). It can be found that the occupied states mainly originate from O-2p and N-2p orbitals, and the unoccupied states are mainly from O-2p, N-2p and P-3p states. The VH process is shown in Fig. S5 in the ESI.† The VH occupied states are mainly contributed by O-2p and N-2p orbitals, and the unoccupied states are mainly contributed by O-2p, N-2p and P-3p states. These results are consistent with the band-resolved analysis (see Fig. 4(b)). It is shown that N-2p and O-2p orbitals make the major contributions to the SHG response in the valence band ($-2.5 \sim 0$ eV) region. In the conduction bands, the P-3p, N-2p and O-2p orbitals in the 6.61–10 eV region make major contributions to the SHG response.

In summary, we suggest the polar $[PN_2O_2]$ tetrahedron as a new FSU for designing non- π -conjugated deep-UV NLO materials because of its large hyperpolarizability, polarizability anisotropy and suitable HOMO–LUMO gap. To confirm this, we performed the crystal structure search for the P–N–O system to explore the potential deep-UV structure. We discovered a thermodynamically stable non-centrosymmetric PNO structure and a competing metastable one. Both of them have a large band gap of 6.61–6.63 eV, large birefringence of 0.151–0.174 @1064 nm

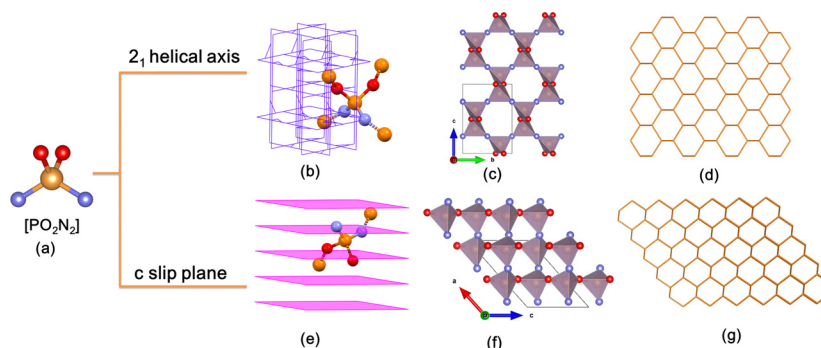


Fig. 3 (a) $[PN_2O_2]$ unit, (b) 2_1 helical axis in *a*, *b*, and *c* directions, (c) 3D structure of **PNO-I**, (d) topology of **PNO-I**, (e) *c* slip plane in the *c* direction, (f) 3D structure of **PNO-II**, and (g) topology of **PNO-II**.

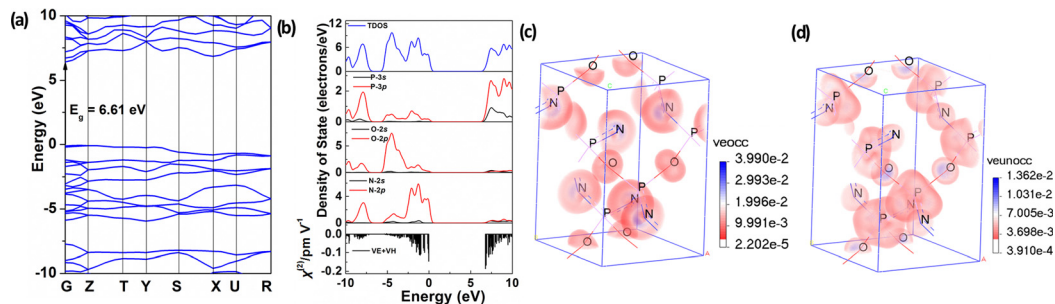


Fig. 4 Band structure (a), density of states band-resolved NLO susceptibility $\chi^{(2)}$ (b), and the SHG-density of occupied (c) and unoccupied (d) states.

and large SHG responses that are about 6~9 times that of KDP. As compared with other deep-UV materials assembled from non- π -conjugated groups and even those π -conjugated groups, PNO has the largest SHG coefficient.

C. W. X. and A. R. O. conceived and designed the research. C. W. X. and A. T. performed all the research work. C. W. X. and A. T. performed the theoretical analysis and wrote the paper. All the authors discussed the results and commented on the manuscript.

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Conflicts of interest

The authors declare no conflict of interest.

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